

Molecular Contamination Modeling with CTSP

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Abstract. Spacecraft instruments and thermal control surfaces are generally highly sensitive to molecular and particulate contamination. Despite best efforts taken during assembly, integration, and test, it is impossible to completely eliminate all sources of contaminants. Contamination transport analysis then becomes of critical importance. It can be used to predict the end of life accumulation on critical surfaces from prescribed source rates. Conversely, given allowable deposition levels, contamination modeling can be used to determine the cleanliness requirements to be met prior to launch. This paper describes a recently developed code for modeling contamination transport. Unlike other tools used by the community, CTSP concurrently traces many simulation particles through small time steps. This allows the code to visualize contaminant partial pressures, and to also include aerodynamic, gravitation, or electrostatic forces. The code is demonstrated by simulating an outgassing characterization test in a bell jar.

INTRODUCTION

Modern spacecraft almost always contain instruments, thermal control surfaces, or mechanical devices that are susceptible to contamination. The term “contamination” refers to the presence of foreign materials that may manifest themselves either as molecular film, or small, dust-like particulates. Molecular films just few monolayers thick can lead to a significant transmission loss in optical instruments. They also lead to degraded performance of thermal control surfaces[1]. Particulates scatter light and can lead to damage of mechanical or high voltage components[2]. Despite best efforts taken during fabrication, integration, and testing, it simply is not possible to completely eliminate all sources of contaminants. Contamination transport modeling then becomes an essential tool in the arsenal of a systems engineer. It can be used to predict the end of life deposition levels given some source rates. Conversely, given the allowable end of life deposition, transport modeling can be used to derive the cleanliness requirements to be met prior to launch.

Historically, different methodologies have been applied to molecular and particulate contaminants. Molecular contaminants, due to their tiny size and mass, are assumed to be unaffected by gravity. They are also assumed to remain neutral. Of interest to the contamination community is the free molecular flow regime in which Knudsen number $Kn \gg 1$. Knudsen number is the ratio of distance traveled by molecules between collisions to a characteristic length, such as the diameter of the vacuum chamber. In free molecular flow, molecules collide with chamber walls (or spacecraft) more frequently than with each other. Intermolecular collisions can be ignored, and molecules can be assumed to travel in straight lines like photons. This formulation leads to the widespread use of radiation heat transfer codes for contamination analysis[3]. The modeling problem reduces to the task of computing gray body view factors between sensitive surfaces and all outgassing sources. Different methods exist for this calculation including radiosity, which is based on matrix inversion, and Monte Carlo ray tracing.

The primary limitation of the radiation-based approach is that by imposing the straight line trajectories, it becomes impossible to consider influence of forces on the contaminants. It is well known that molecules exposed to solar radiation undergo photoionization and subsequently backflow to the spacecraft[4]. In the larger particulates, gravitational and aerodynamic forces become non-negligible. Historically this required development of specialized codes for each contributing factor. Some examples include the electrostatic return ESR code[5] and Mastram’s particle module for analyzing redistribution in a launch vehicle fairing[6]. A new simulation code CTSP (Contamination Transport Simulation Program) was developed to address these shortcomings. This code is described in this paper. The paper begins by introducing CTSP and describing the surface model used in the molecular transport analysis. The code is

then demonstrated by simulating an experimental method for determining outgassing rate of a test article placed in a vacuum chamber.

CODE OVERVIEW

Particle Push

CTSP follows methodology common in Particle in Cell (PIC)[7] or Direct Simulation Monte Carlo (DSMC)[8] codes, and concurrently traces many simulation *particles* through small Δt time steps. With the exception of surface interaction handling, CTSP does not distinguish between molecular and particulate contaminants. Particulates are simply simulation particles with larger mass and larger cross-sectional surface area. By pushing particles through small time steps, it is possible to consider velocity change due to action of forces. At each time step, velocities are updated from Newton's second law, $d\vec{v}/dt = \vec{F}/m$ where \vec{v} is the velocity, m is the particle mass, and \vec{F} is the summation of all forces acting on the particle. CTSP supports gravitational, electrostatic, aerodynamic, and solar pressure forces. Orbital motion of the parent body can also be defined to model particulate return on orbit crossings. Particle position is advanced by integrating $d\vec{x}/dt = \vec{v}$ using the Leapfrog method. The concurrent push of many particles introduces two benefits. First, it is now possible to visualize gas bulk properties such as contaminant partial pressure. Secondly, although not yet implemented, the code could be fitted with a DSMC subroutine to consider cases outside the realm of free molecular flow. One application of interest in this transition region is modeling of chamber repressurization.

Geometry and Simulation Inputs

Simulation geometry is described by a mixed triangular/quadrilateral surface mesh. This mesh is typically generated from a CAD drawing using finite element analysis meshing programs. Surface elements corresponding to the same logical component need to be collected into groups. The grouping is important since CTSP assigns surface properties only at the group level. The groups are also used to specify particle sources. The code implements multiple injection models applicable to contamination modeling, including outgassing (described below), effusion to model vents, particulate distribution according to the IEST-STD-1246 standard[9], and random coverage by fibers. CTSP does not contain built-in field solvers, and the appropriate fields needed to compute forces need to be provided by the user. For instance, in order to model particulate fallout in a moving air field, it is necessary to run a CFD analysis to generate a spatial distribution of gas velocity vectors and density scalars. These quantities are then interpolated onto the particle position. PIC and DSMC codes generally utilize a volume mesh to perform this interpolation. In developing CTSP, decision was made to avoid the need for a volume mesh. In the flow regimes considered by CTSP, there is no appropriate physical parameter usable for sizing the volume cells. In PIC, the cell size is typically based on the local Debye length. In DSMC, the collision mean free path dictates the cell size. But in free molecular flows, the cell size could be chosen arbitrarily. Typical CFD solution will contain near-wall region in which the velocity changes rapidly. Interpolating these values onto a too coarse Cartesian grid would result in loss of these fine features. On the other hand, utilizing the finite element grids from the field solvers would also be impractical, as it would require the code to track which cell each particle is located for each input grid.

Octree and Surface Interactions

Hence, an octree is used to efficiently store the point cloud of field quantities. Prior to each particle push, CTSP first defines a control box about the particle's position. An initially small box is expanded until enough points are retrieved. The located data points $u_i(\vec{x}_i)$ are interpolated onto the particle position \vec{x} using the inverse distance weighting method, $u(\vec{x}) = \sum_{i=1}^N w_i(\vec{x})u_i / \sum_{i=1}^N w_i(\vec{x})$ with $w_i(\vec{x}) = d(\vec{x}, \vec{x}_i)^{-p}$. Here d is the distance between the points, and p is the interpolation smoothing factor. An octree is also used to store the surface mesh elements. The starting and final position of a particle during a single push is used to define a bounding box. All surface elements located within the bounding box are retrieved from the octree. The code checks for impact by first performing line-plane intersections of the line segment traversed by the particle with the plane of the surface element. The intersections are sorted by the distance from the particle starting point. This step is necessary to account for particles completely traversing through a thin surface in a single time step. Next, the intersection point is tested for inclusion within the surface element triangle or quadrangle by looping over all vertices, and forming bisections of vertex angle with a segment to the intersection point. The point is located within the enclosed polyhedral if $\angle_1 + \angle_2 = \angle_0$ for all vertexes. If intersection is found,

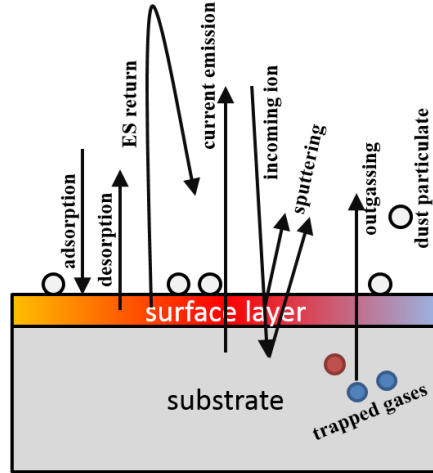


FIGURE 1. Detailed surface interaction model. Only outgassing and surface adsorption/desorption are considered in this work.

the particle is first moved to the surface. The code then calls an intersection handler. This is one areas where CTSP distinguishes between particulates and molecules. If the impacting particle is a particulate, the intersection handler computes the post-impact velocity using a user defined coefficient of restitution, $v_2 = \alpha_r v_1$. If the resulting velocity magnitude is less than some threshold, the particle is assumed to become stuck. Its dimensions are then used to update a histogram of surface particulate distribution. Otherwise, the particle is given a new random velocity away from the surface and the particle is pushed for the remainder of the time step. Behavior of impacting molecules is based on the surface residence time, as described below.

Outgassing and Desorption

As noted in introduction, it is customary to analyze molecular transport by computing gray body view factors between sources and sensitive surfaces. The view factor is then multiplied by the source outgassing rate to obtain the contribution reaching the target. Besides imposing the straight line trajectory, this approach also ignores the surface-vacuum interface. Molecular outgassing arises from volatile gases trapped inside the material diffusing to the surface and desorbing into gas phase. The outgassing rate is not constant. As the amount of trapped material gets depleted, the outgassing rate also goes down. The rate also scales with temperature. It is in fact common to perform a vacuum bakeout of spacecraft components to reduce their initial outgassing rate. The model in CTSP attempts to capture this temperature dependence. The model is based on the conceptual view shown in Figure 1. All simulation surfaces are assumed to consist of a native substrate containing an arbitrary heterogeneous mixture of trapped gases. At the surface is a thin surface layer also composed of an arbitrary combination of molecular species. Initial compositions of the two regions is specified in the input file for each surface element group as illustrated below:

```
comp{name:ebox, mat:a1, trapped_mass:1e-4, trapped_mats:[0.85*hc1,0.10*water,0.05*n2],
      surf_h:5e-10, surf_mat:hc1, temp:350}
```

The amount of mass lost to diffusion is given by[1]

$$\frac{dm}{dt} = Cm \exp(-E_a/RT)^r \quad (1)$$

where C is a reaction constant, m is the mass, E_a is the material activation energy (typically in kcal/mol), R is the gas constant (kcal K/mole), and T is the temperature (K). For diffusion-limited processes, $r = -0.5$, although the example presented in this paper used $r = 0$ (no time dependence). The reaction constant and activation energy can be determined experimentally with the industry standard ASTM-E1559 test[10]. E_a can also be obtained by performing a thermogravimetric analysis (TGA) with a quartz crystal microbalance (QCM). In the absence of either of these data sets, we can estimate activation energy by considering molecular residence time. On a low velocity surface impact, the

molecule temporarily adheres to the surface and establishes thermal equilibrium. The molecule remains attached until, due to random events, it acquires enough energy to escape the electrostatic attraction to the surface. The molecule then leaves in direction following the cosine law from the surface normal[11]. The time spent on the surface is known as residence time,

$$\tau_r = \tau_0 \exp(E_a/RT) \quad (2)$$

Here τ_0 is the oscillation period of the molecule with typical value $\sim 10^{-13}$ s[1]. For most species, the residence time will be very short except on surfaces below that gas material's condensation temperature. We can use activation energies to generate materials characteristic of species commonly encountered in thermal vacuum tests. For instance, molecule with $E_a = 12$ kcal/mol will have τ_r of less than 10^{-5} s on a 350K surface. The residence time increases to over 10,000 s if the surface is cooled to 150K. This value of E_a can thus be used to characterize a generic hydrocarbon that is driven off at warmer temperatures but sticks to all cold surfaces. A contaminant layer will build up on the surface as long as the arrival rate exceeds the rate of departure. For deposition thickness less than a monolayer, the rate of evaporation is given by first order desorption[11],

$$\frac{dN}{dt} = \frac{1}{\tau_r} N \quad (3)$$

Numerical Implementation

When a molecule impacts a surface, CTSP computes residence time using equation 2 at the temperature of the impacted surface element. If $\tau_r \leq \Delta t$, the particle is re-emitted. First new velocity magnitude v_r is sampled from the Maxwellian speed distribution function at the surface temperature. The particle's new speed is set to $v + \alpha_a(v_r - v)$, where α_a is the thermal accommodation coefficient. Full accommodations $\alpha_a = 1$ was assumed in the example presented in this paper. The new velocity direction follows the cosine law, which is sampled according to [8]. On the other hand, if $\tau_r > \Delta t$, the particle is adsorbed into the surface layer. Numerically, this involves removing the particle from the simulation and incrementing surface count of particle's material i by the specific weight, $\Delta N_i = w_{sp}$. The counts can be related to surface layer thickness by $h = N(4/3)\pi r^3/A_{el}$, where r is the molecular radius. Outgassing is implemented as a source called by the main loop at every time step. The algorithm computes the amount of outgassed mass for each component from equation 1. The mass, converted to number of molecules, is moved to the surface layer. The incremental value for each surface element j is scaled by the local area ratio, A_j/A_{tot} , where A_{tot} is the total surface area of the component. Next, the code loops through all surface elements, and for each computes the amount of mass lost to desorption per equation 3. For cold surfaces, this value will be too small to produce particles. The surface layer is depleted according to the number of generated particles. The benefit of this approach is that instead of utilizing surface sticking coefficients, the impact behavior is determined solely by material properties and surface temperature. The overall algorithm can be summarized by the pseudocode below:

```

load material data (density, molecular radius, TML, ...)
load component data (temperature, total mass, cleanliness level, ...)
initialize surface data structures

for each timestep:
  #particle push
  for each particle:
    update velocity and position
    if surface impact:
      compute residence time
      if greater than time step:
        deposit to surface layer
      else:
        re-emit following cosine law

  #outgassing model
  for each component:
    compute outgassing mass loss
    add to surface layer

```

```

#desorption
for each surface element:
  compute desorption loss
  inject particles with cosine distribution
  deplete surface layer

```

NUMERICAL EXAMPLE

To demonstrate the code, we consider a commonly encountered engineering task: characterizing the outgassing rate of a test article exposed to vacuum. Figure 2 shows the setup. The test article (a harness) is placed on a platen inside a bell jar. A curved duct connects the bell jar to a vacuum pump. Inside the bell jar is a scavenger plate that can be flooded with cryogenic nitrogen to create a secondary pumping surface. Single quartz crystal microbalance (QCM) is also placed near the harness. This test corresponds to a “hot-wall” configuration in which the bell jar surface and the platten are not cold enough to condense the contaminant.

View factor

QCMs indirectly measure mass deposition rate on the exposed crystal. However of interest is the outgassing rate from the test hardware. We can start with a mass conservation equation,

$$\Gamma_h A_h - \Gamma_q A_q - \Gamma_p A_p = 0 \quad (4)$$

where Γ and A correspond to mass flux and surface area, respectively, and the subscripts h , q , and p stand for hardware, QCM, and pump. The pump term is a placeholder for all other surfaces cold enough to collect the contaminant. If the QCM and pumps are equally cold sinks, $\Gamma_q = \Gamma_p$, and the above equation simplifies to $\Gamma_h A_h k_{h \rightarrow q} = \Gamma_q A_q$ where

$$k_{h \rightarrow q} = \frac{A_q}{A_q + A_p} \quad (5)$$

is the view factor. In other words, knowing $k_{h \rightarrow q}$, we can relate the deposition rate on the QCM to the outgassing rate of the hardware. Although this calculation appears trivial, the difficulty arises from not having a value for A_p , the effective pump area. Simply considering the cross-sectional area of the pump would ignore pumping inefficiencies and, more importantly, would ignore conductance losses through the duct[11]. Fortunately, we can obtain the pump area by temporarily introducing a secondary collecting surface of a known area[12]. The scavenger plate in the model serves this role. We can then write two equations, one without and one with the scavenger plate,

$$\Gamma_h A_h \left(\frac{A_q}{A_q + A_p} \right) = (\Gamma_q)_1 A_q \quad (6)$$

$$\Gamma_h A_h \left(\frac{A_q}{A_q + A_p + A_s} \right) = (\Gamma_q)_2 A_q \quad (7)$$

The hardware outgassing term Γ_h can be eliminated by substitution giving us

$$A_p \left[(\Gamma_q)_1 - (\Gamma_q)_2 \right] = A_q \left[(\Gamma_q)_2 - (\Gamma_q)_1 \right] + A_s (\Gamma_q)_2 \quad (8)$$

Knowing the scavenger plate area and collecting QCM deposition rates without $(\Gamma_q)_1$ and with $(\Gamma_q)_2$ the scavenger plate, we can use the above equation to solve for the effective pump area. This value is then substituted into Equation 5 to obtain the hardware outgassing rate.

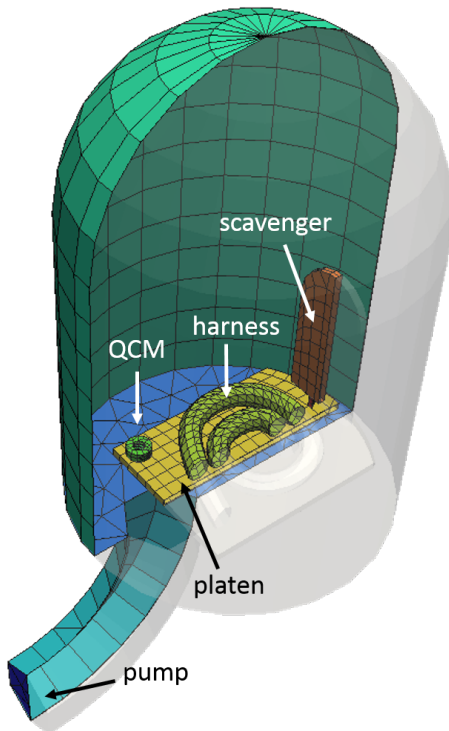


FIGURE 2. Simulation geometry with the major components identified

RESULTS

The simulation was run for 60,000 time steps at $\Delta t = 10^{-4}$ s. The harness initially contained 0.1 g of a trapped hydrocarbon with $E_a = 12$ kcal/mol and outgassing rate constant $C = 1$ s $^{-1}$. The harness temperature was kept at 350K. The pump temperature was set to 15K and the QCM crystal was set to 220K. All other surfaces, including the scavenger, were initially at 300K. At time step 30,000, the scavenger was “flooded” by decreasing its temperature to 100K. The simulation contained approximately 70,000 particles at steady state and warm scavenger. The count decreased to 4,500 once the scavenger was activated due to the increase in total pumping area. We can see the corresponding pressure drop in Figure 3. This Figure plots the partial pressure of the contaminant as well as the surface deposition thickness. The partial pressure was computed by interpolating particles onto a three-dimensional Cartesian mesh. As noted earlier, this mesh is used only for visualization and is not used during the particle push. Particle positions and velocities were interpolated to compute density and temperature, respectively. The partial pressure was computed from $P = nkT$. The interpolation took place every 5 time steps and was used to update a running average. The volume data was saved to a file every 200 time steps for animation. The cumulative data were cleared after each export. The first plot in Figure 3(a) shows the steady-state partial pressure of the contaminant with the warm scavenger. Surface thickness of the deposited molecular film is also shown. Only the QCM crystal and the pump are cold enough to collect contaminants and hence are the only surfaces with non-zero deposition thickness. The situation after the scavenger is activated is plotted in Figure 3(b). Chamber pressure decreased by approximately one order of magnitude from 5×10^{-6} Torr to 5×10^{-7} Torr. Non-zero deposition is now seen on the scavenger plate.

Pump Area

Mass deposited on the QCM crystal versus time step is shown in Figure 4. Over the short duration of the simulation, the hardware outgassing rate is expected to appear constant. This is confirmed by the linear rate of mass increase. We can also see a clear difference in slope before and after the scavenger activation at time step 30,000. Microsoft Excel’s trendline was used to fit linear curves to the two segments. Of interest is the ratio of slopes. The effective pump area

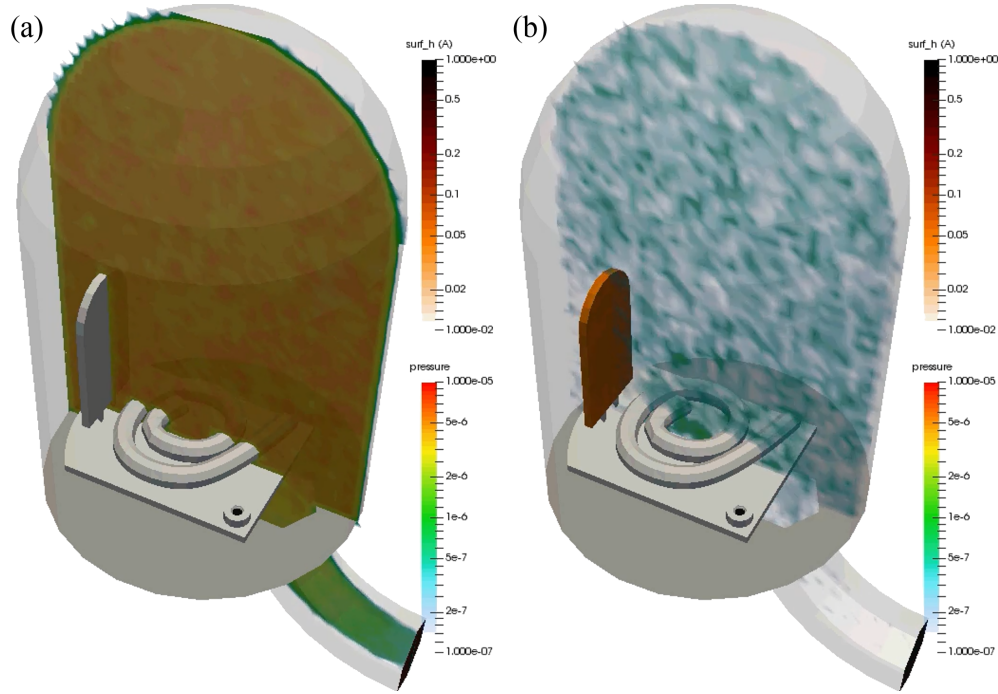


FIGURE 3. Contaminant partial pressure before (a) and after (b) scavenger activation

is calculated from Equation 8 by substituting these two slopes for $(\Gamma_q)_1$ and $(\Gamma_q)_2$. Areas of the QCM crystal, A_q , and the scavenger plate, A_s , were obtained by summing the appropriate triangular and quadrilateral mesh elements. The effective pump area obtained from the simulation is 0.00278 m^2 . This corresponds to reduction by ratio of 0.429 from the geometric area of 0.00649 m^2 .

Outgassing Rate

We can now compute the harness outgassing rate,

$$r_h = r_q \frac{1}{k_{h \rightarrow q}} \frac{A_q}{A_h} \quad (9)$$

where, for the warm scavenger case,

$$k_{h \rightarrow q} = \frac{A_q}{A_q + A_p} = 0.0353 \quad (10)$$

In the above equation, $r_0 \equiv \Gamma_0 A_0$. Hardware outgassing rate, as measured by the QCM, is $r_h = 2.85 \times 10^{-12} \text{ kg s}^{-1}$. This value can be compared against the expected value from the source model. The production rate of the contaminant is given by Equation 1 and evaluates to $3.214 \times 10^{-12} \text{ kg s}^{-1}$. The 11% error could likely be reduced by increasing the number of simulation particles.

CONCLUSION AND FUTURE WORK

New simulation program was developed to model transport of molecular and particulate contaminants. Unlike other tools used by the community, CTSP concurrently traces multiple particles through small simulation time steps. At each time step, particle velocities are updated based on user-provided force fields. The concurrent push also allows the code to visualize contaminant partial pressures. CTSP implements several source models applicable to contamination modeling, including one for molecular outgassing and adsorption. This model was demonstrated by simulating a

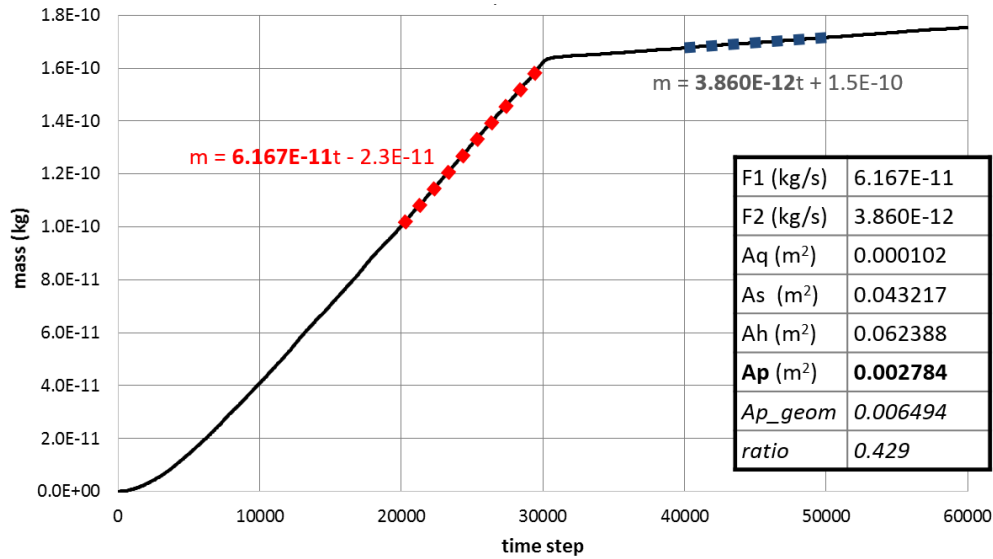


FIGURE 4. Mass collected by the QCM. Difference in rates before and after scavenger plate activation is clearly seen. Computed pump area is also shown.

common engineering task of characterizing the outgassing rate of a test article. A simulation QCM was used to measure deposition rate. Taking into account the effective pump area, it is possible to backtrack the test article outgassing rate from the QCM deposition data. The obtained result agreed within 11% from the expected value. Future work includes additional code validation, ideally using experimental data. Furthermore, code parallelization is currently in progress, and preliminary work has been completed on a DSMC module for analyzing the transition region.

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